# An analytical method for estimation of parameters of Self Mixing Interferometric phase equation over all feedback regimes

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#### Abstract

This paper presents a novel algorithm for measuring the linewidth enhancement factor of semiconductor lasers and the optical feedback level factor in a semiconductor laser with an external cavity. The proposed approach is based on analysis of the self-mixing phase equation to deduce equations for finding parameters given only knowledge of the perturbed phase. The effectiveness of the method has been validated with accuracy of 8.6% and 1.7% for 'C' and alpha respectively while covering all feedback regimes.

# An analytical method for estimation of parameters of Self Mixing Interferometric phase equation over all feedback regimes

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8 Abstract. This paper presents a novel algorithm for measuring the linewidth enhancement 9 factor of semiconductor lasers and the optical feedback level factor in a semiconductor laser 10 with an external cavity. The proposed approach is based on analysis of the self-mixing phase 11 equation to deduce equations for finding parameters given only knowledge of the perturbed 12 phase. The effectiveness of the method has been validated with accuracy of 8.6% and 1.7% for 13 C and  $\alpha$  respectively while covering all feedback regimes.

# 14 1. Keywords

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Linewidth enhancement factor (LEF), optical feedback, self-mixing interferometry, semiconduc tor lasers

## 17 2. Introduction

In optical feedback self-mixing interferometry (OFSMI), the instantaneous distance between 18 the laser semiconductor diode driven by a constant injection current and a remote surface 19 which back-scatters a small amount of optical power back into the laser diode cavity[7]. 20 Linewidth enhancement factor (LEF) (denoted as  $\alpha$ ) is a fundamental parameter of self-mixing 21 interferometry as it characterizes the linewidth, the chirp, the injection lock range, and the 22 response to optical feedback [5]. The influence of the parameters C and  $\alpha$  on the emitted laser 23 intensity has been extensively analyzed in [1] and [2]. Establishing an accurate measurement has 24 been a challenging and active research topic that has attracted extensive research work during 25 the past two decades [5]. Existing approaches include the direct physical measurement of the 26 subthreshold optical spectrum as the injected current is varied [3] and techniques based on the 27 analysis of the locking regimes induced by optical injection from a master laser [4]. Moreover, 28 an analytic method, based on gradient descent approach was presented [6], which showed the 29 accuracy of 6.7% and 4.63% for C and respectively. 30

## 31 3. Main theory

When the optical feedback phenomenon occurs, the laser wavelength is no longer the constant  $\lambda_o$  but is slightly modified and becomes a function of time  $\lambda_f(t)$  when D(t) varies. The wavelength

fluctuations can be found by solving the phase equation,

$$x_o(t) = x_f(t) + C\sin(x_f(t) + \arctan(\alpha)) \tag{1}$$

 $x_o$  and  $x_f$  are referred as perturbed and unperturbed phase respectively,  $\alpha$  is the linewidth enhancement factor and C is the coupling factor.  $x_o$  and  $x_f$  can be represented as,

$$x_o(t) = 2\pi\nu_o(t)\tau(t) \tag{2}$$

$$x_f(t) = 2\pi\nu_f(t)\tau(t) \tag{3}$$

where  $\tau(t) = 2D(t)/c$  is the round trip time, with c as speed of light.  $\nu_f(t)$  and  $\nu_o(t)$  represents optical frequencies with and without feedback. Laser feedback output optical power(LDOOP) P(t) depends on the SM phenomenon and written as:

$$P(t) = P_o[1 + m\cos(x_f(t))]$$

where  $P_o$  is the power emitted by the free-running state laser diode and m is the modulation index. Therefore, for purpose of displacement measurement, we track from SM signal P(t)

<sup>33</sup> index. Therefore, for purpose of displacement measurement, we track from SM signal P(t)<sup>34</sup> measurement toward perturbed phase  $x_f(t)$  to  $x_o(t)$  toward displacement D(t) measurement.Due

 $_{35}$  to environmental fluctuations and uncertainty, it is desirable to estimate parameters C and  $\alpha$ 

<sup>36</sup> for a given condition of the interferometer for robust displacement measurement.

#### 37 4. Theoretical derivation

38 4.1. Estimation of alpha

Taking derivative on both sides of (1),

$$\frac{dx_o}{dt} = \frac{dx_f}{dt} (1 + C\cos(x_f + \arctan(\alpha)))$$
(4)

$$\frac{\frac{dx_o}{dt}}{\frac{dx_f}{dt}} = 1 + C\cos(x_f + \arctan(\alpha))$$
(5)

When extremas of  $x_o$  and  $x_f$  are reached simultaneously, then their derivatives approach zero simultaneously and their ratio approaches 1 by taking limit on time, i.e., when  $t = t_{ext}$ ,

$$\lim_{t \to t_{ext}} \frac{\frac{dx_o}{dt}}{\frac{dx_f}{dt}} = 1 \tag{6}$$

Considering when  $t = t_{ext}$ , then  $x_f = x_{fe}$  and  $x_o = x_{oe}$ , Putting in (5),

$$\cos(x_{fe} + \arctan(\alpha)) = 0 \tag{7}$$

$$x_{fe} + \arctan(\alpha) = k\pi - \frac{\pi}{2} \tag{8}$$

Putting equation(8) in equation(1), as  $\sin(k\pi - \frac{\pi}{2}) = \pm 1$ 

$$x_{oe} = x_{fe} \pm C \tag{9}$$

From equation(1),  $x_o - x_f$  is maximally bounded by constant C, when  $x_f$  is allowed to increase from 0 to some value (less than  $\frac{\pi}{2} - \arctan(\alpha)$ ). Therefore, on local maximas of displacement,  $x_o$ swings  $x_f$  by C. Lets denote maximas of  $x_o$  and  $x_f$  as  $x_{om}$  and  $x_{fm}$  respectively, then on local maximas of displacement,

$$\begin{aligned} x_{om} &= x_{fm} + C \\ \text{Page 2 of 8} \end{aligned} \tag{10}$$

Putting in equation(1),

 $x_{fm} + C = x_{fm} + C\sin(x_{fm} + \arctan(\alpha)) \tag{11}$ 

$$\sin(x_{fm} + \arctan(\alpha)) = 1 \tag{12}$$

$$x_{fm} + \arctan(\alpha) = \frac{\pi}{2} + 2k\pi \tag{13}$$

where, **k** is an integer.

$$\alpha = \tan\left(\frac{\pi}{2} + 2k\pi - x_{fm}\right) \tag{14}$$

Due to periodicity of tan function by  $2k\pi$ , equation (12) becomes,

$$\alpha = \tan(\frac{\pi}{2} - x_{fm}) \tag{15}$$

$$\alpha = \tan(\arcsin(1) - x_{fm}) \tag{16}$$

39 4.2. Estimation of C

Now, the perturbed phase  $(x_f)$  being modified form of  $x_o$  exhibits sharp transitions (from moderate to high feedback regime but it can be extended to all regimes once the algorithm is developed), then at those transitions derivative of  $x_f$  approaches infinity. Let  $x_f$  be represented as  $x_{ft}$  on transitions of perturbed phase, then equation (5) becomes,

$$1 + C\cos(x_{ft} + \arctan(\alpha)) = 0 \tag{17}$$

$$C = \frac{-1}{\cos(x_{ft} + \arctan(\alpha))} \tag{18}$$

## 40 5. Algorithm Design

<sup>41</sup> 5.1. Algorithm for alpha estimation Since derivative of  $\arcsin(t)$  is given by,

$$\frac{d(\arcsin(t))}{dt} = \frac{1}{\sqrt{1-t^2}} \tag{19}$$

Around t=1,  $\frac{\operatorname{darcsin}(t)}{dt}$  becomes quite large and sensitive, and slight deviations can lead to huge errors. Therefore, we can add counter sensitivity term as  $\epsilon$  which has following properties.

$$\epsilon \ll 1 \text{ and } \epsilon \ge 0$$
 (20)

So, equation(14) becomes modified as,

$$\alpha = \tan(\arcsin(1 - \epsilon) - x_{fm}) \tag{21}$$

<sup>42</sup> If we are given information that  $\alpha$  is bounded by some threshold ' $\tau \ge 0$ ' and 0, then we find  $\alpha$ <sup>43</sup> by following algorithm

Step 1: Initially choose  $\epsilon = 0$  and step size of varying it as  $\delta \epsilon$ 

Step 2: For a given sample of perturbed phase  $x_f$ , find the maximum positive element as  $x_{fm}$ , which would be global maxima of  $x_f(t)$ .

Step 3: Plug the values into equation(21). If  $\alpha$  is outside  $[0, \tau]$ , then increase  $\epsilon$  to  $\epsilon + \delta \epsilon$  and repeat 48 Step 3. Otherwise,  $\alpha$  would be required estimation.

- 49 5.2. Algorithm for C estimation
- <sup>50</sup> Due to sharpness of SM signal, a slight deviation of arguments of equation(18), can create huge
- <sup>51</sup> errors. Due to this non-smooth property, we have to use statistic of finding globally sharpest
- <sup>52</sup> transition and exploiting equation (18).
- Step 1: For suitable magnitude threshold M, calculate derivative of sample of  $x_f$  as  $\frac{dx_f}{dt}$  and find the set X as,

$$X = \{x_f(t)|t = \arg[\frac{dx_f}{dt} \ge M]\}$$
(22)

Step 2: Plug the values into equation(18), then estimated value of C is,

$$C = max(\frac{-1}{\cos(x + \arctan(\alpha))}), \,\forall x \in X$$
(23)

# 53 6. Simulation Results

54 6.1. Estimation of C and alpha

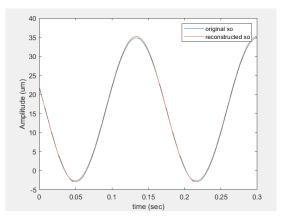


Figure 1. Reference unperturbed phase and reconstructed unperturbed phase for a displacement of 6kHz

6.1.1. Remark The above-mentioned algorithm can be extended and conjugated with an
 optimization algorithm, as depicted in [6], which would decrease the number of iterations
 for convergence and increase accuracy. Moreover, an obvious observation is that an arbitrary
 smoothing digital filter could increase accuracy when reconstructing displacement.

#### 59 6.2. Reconstruction of unperturbed phase and displacement

For a sinusoidal displacement of 3 peak to peak  $\mu$ m and 6kHz, sampled at 10kHz with reference values of C and  $\alpha$  as 4 and 5 respectively, on applying mentioned algorithms with  $\tau = 0.02$ and M = 10000, we get a reconstructed unperturbed phase which gives mean square error of 32.4nm, by utilizing equation(1), as shown in figure(2) and (3) and reconstructed displacement

- <sup>64</sup> is mentioned in figure(4). Displacement is reconstructed by using equation (2).
- $_{65}$  Similarly, for a sinusoidal displacement of 3 peak to peak  $\mu$ m and 6kHz,12kHz and 18kHz,
- sampled at 10kHz with reference values of C and  $\alpha$  as 4 and 5 respectively, on applying mentioned
- algorithms with  $\tau = 0.19$  and M = 10000, we get a mean square error of 96.1nm and results
- are mentioned in figure (5), (6) and (7), after utilizing equation (1) and (2). To work with the
- <sup>69</sup> mentioned method, the "feedback regimes" term would be ignored and values of C and  $\alpha$  would <sup>70</sup> be randomly experimented with, from 1 to 9 for C and 2 to 5 for  $\alpha$  respectively. For simulation,
- <sup>71</sup> a sample of the perturbed phase corresponding to a sinusoidal displacement of 3 peak to peak

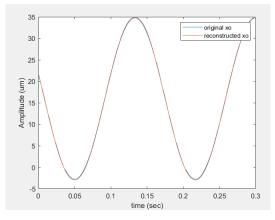


Figure 2. Plot of reference unperturbed phase and reconstructed unperturbed phase based on estimated C and  $\alpha$  for a displacement of 6kHz

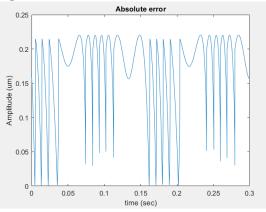


Figure 3. Plot of absolute error between reference unperturbed phase and reconstructed unperturbed phase based on estimated C and  $\alpha$  for a displacement of 6kHz

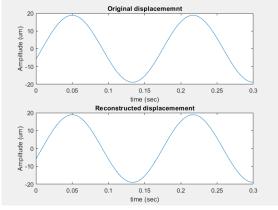


Figure 4. Reference displacement and reconstructed displacement for a displacement of 6kHz

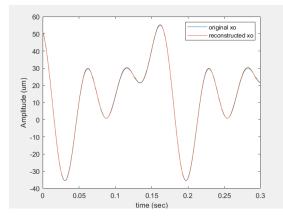


Figure 5. Plot of reference unperturbed phase and reconstructed unperturbed phase based on estimated C and  $\alpha$  for a displacement having 6kHz,12kHz and 18kHz components

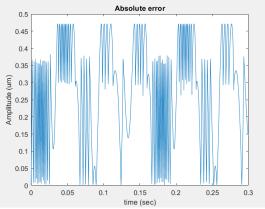
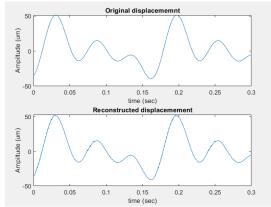


Figure 6. Plot of absolute error between reference unperturbed phase and reconstructed unperturbed phase based on estimated C and  $\alpha$  for a displacement having 6kHz,12kHz and 18kHz components



**Figure 7.** Reference displacement and reconstructed displacement for a displacement having 6kHz,12kHz and 18kHz components

 $\mu$ m and 6kHz, sampled at 10kHz, was used. The following table represents values of estimated

values of C corresponding to reference C values, with fixed  $\alpha = 5$ , along with  $\tau = 0.02$  and

74 M = 10000 found by hit and trial.

Estimated $C(\alpha = 5)$	Reference $C(\alpha = 5)$	Percentage error $\%$
1.1708	1	17
2.3401	2	17
4.0779	4	1.9475
5.4588	5	9.17
6.6384	7	5.7
9.1011	9	1.11

<sup>76</sup> The average error for above table is 8.67%, for estimating C independently.

 $_{77}$   $\,$  Similarly, following table presents estimated values of  $\alpha$  corrosponding to reference values, when

 $_{78}$  C is kept constant at 5  $\,$ 

Estimated $\alpha(C=5)$	Reference $\alpha(C=5)$	Percentage error $\%$
2.1098	2	5.49
3.0080	3	0.2667
4.0013	4	0.0325
5.0620	5	1.24

<sup>80</sup> The average error for above table is 1.75%, for estimating  $\alpha$  independently.

Simultaneous prediction of values of C and  $\alpha$  corresponding to reference values are presented in following table.

The average error for estimating C and  $\alpha$  simultaneously is 3.07% and 22.35% respectively. It should be noted that the above-presented data covers all feedback regimes.

Even though some errors appear large, for example, for reference C=6,  $\alpha$ =4, give percentage

errors of 3.16% and 32.25% respectively, when used for reconstruction of unperturbed phase from equation(1) gives means square error of 100nm, which is quite reasonable, and its depiction is

 $^{89}$  presented in figure(1).

# 90 7. Conclusion

<sup>91</sup> Therefore, based on SMI phase equation (1), the algorithm has been developed for estimation of <sup>92</sup> C and  $\alpha$  and is useful for reconstruction of unperturbed phase from sole knowledge of perturbed <sup>93</sup> phase, for measurement of displacement of the target, under situations where C and  $\alpha$  are known <sup>94</sup> or vary with the accuracy of 8.67% and 1.75% respectively, when independent estimation and <sup>95</sup> 22.35% and 3.07% for simultaneous estimation. But even still, the mean square error bound for

<sup>96</sup> reconstructed displacement is less than 100nm.

# 97 8. Acknowledgment

98 This study was not funded.

# 99 9. Conflict of Interest

<sup>100</sup> The author declares that he has no conflict of interest.

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#### 101 **10. References**

- [1] G Acket, Daan Lenstra, A Den Boef, and B Verbeek. The influence of feedback intensity on longitudinal mode
   properties and optical noise in index-guided semiconductor lasers. *IEEE Journal of Quantum Electronics*,
   20(10):1163–1169, 1984.
- [2] Silvano Donati, Guido Giuliani, and Sabina Merlo. Laser diode feedback interferometer for measurement of
   displacements without ambiguity. *IEEE journal of quantum electronics*, 31(1):113–119, 1995.
- [3] Piotr Konrad Kondratko, Shun-Lien Chuang, Gabriel Walter, Theodore Chung, and Nick Holonyak Jr.
   Observations of near-zero linewidth enhancement factor in a quantum-well coupled quantum-dot laser.
   Applied physics letters, 83(23):4818-4820, 2003.
- [4] G Liu, Xiaomin Jin, and Shun-Lien Chuang. Measurement of linewidth enhancement factor of semiconductor
   lasers using an injection-locking technique. *IEEE photonics Technology letters*, 13(5):430–432, 2001.
- [5] Marek Osinski and Jens Buus. Linewidth broadening factor in semiconductor lasers-an overview. *IEEE Journal of Quantum Electronics*, 23(1):9–29, 1987.
- [6] Jiangtao Xi, Yanguang Yu, Joe F Chicharo, and Thierry Bosch. Estimating the parameters of semiconductor
   lasers based on weak optical feedback self-mixing interferometry. *IEEE Journal of Quantum Electronics*,
   41(8):1058–1064, 2005.
- 117 [7] Yanguang Yu, Yuanlong Fan, and Bin Liu. Self-mixing interferometry and its applications. In Optical Design
- and Testing VII, volume 10021, page 100210U. International Society for Optics and Photonics, 2016.